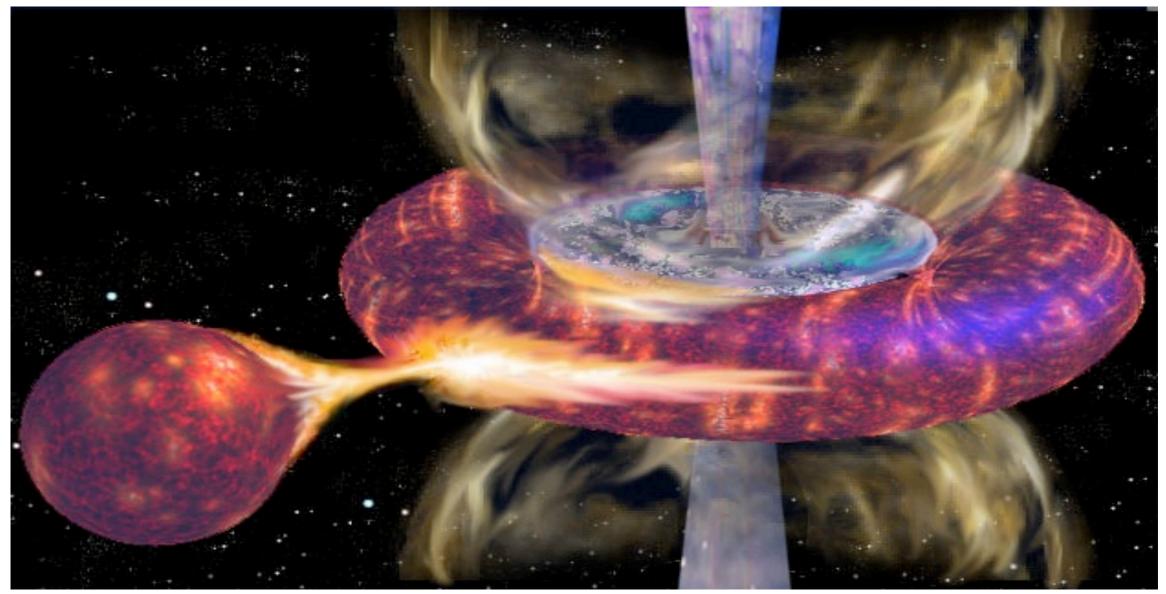


# A disk instability based classification for microquasars

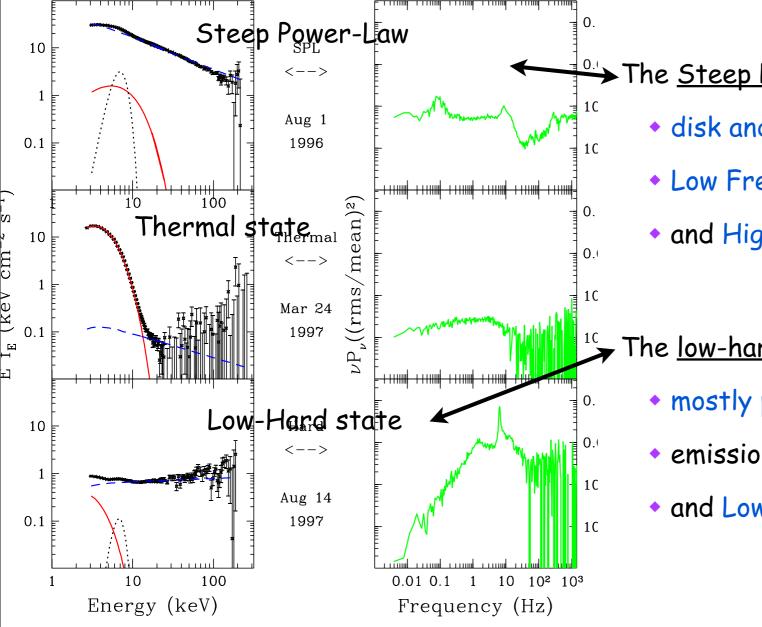


#### P. Varniere, M. Tagger, J. Rodriguez

#### Microquasar observations

a microquasar is a binary system hosting a black-hole of a few tens of solar masses being "fed" an accretion disk and emitting relativistic jets.

Microquasars appear to transit between several states, here we are m ones displaying timing variability known as Quasi-Periodic Oscillations GRO J1655-40



<sup>⊙.</sup>→The <u>Steep Power-Law state</u> involves

- disk and power-law emissions are both present
- Low Frequency Quasi Periodic Oscillation (LFQPO)
- and High Frequency Quasi Periodic Oscillation (HFQPO)

The <u>low-hard state</u> involves

- mostly power-law spectrum
- emission of a steady jet
- and Low Frequency Quasi Periodic Oscillation

large variety of observations (X-rays, Gamma rays, Infra-Red, radio,...) and techniques (timing, lags)

→ many constraints on theories

### AEI and LFQPO

in a disk threaded by a vertical B  $\sim$ 

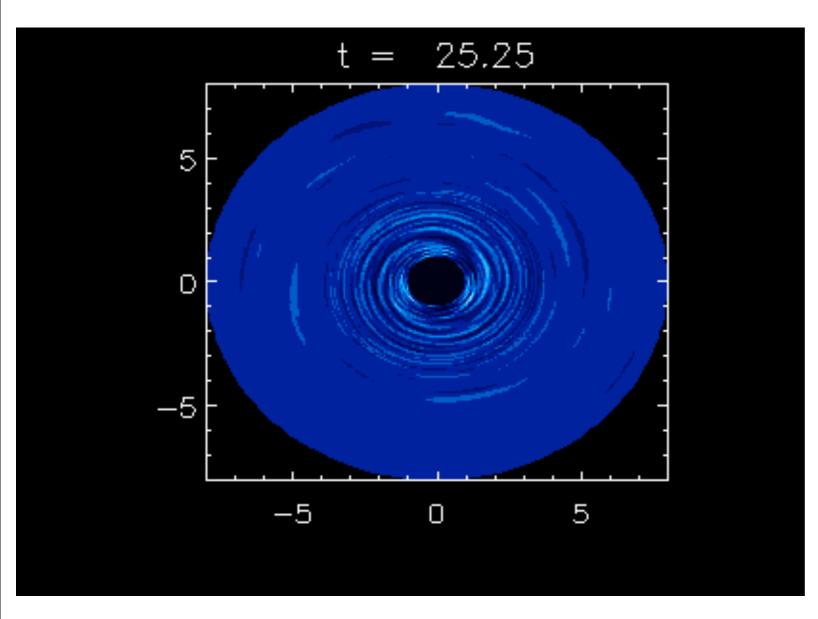
equipartition

needs positive gradient of magnetovortensity:

$${\cal L}_B \;=\; {\kappa^2\over 2\Omega}\; {\Sigma\over B^2}$$

→ application to LF-QPO: frequency/radius, correlation with jet/radio, global properties...

# AEI and LFQPO



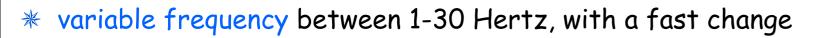
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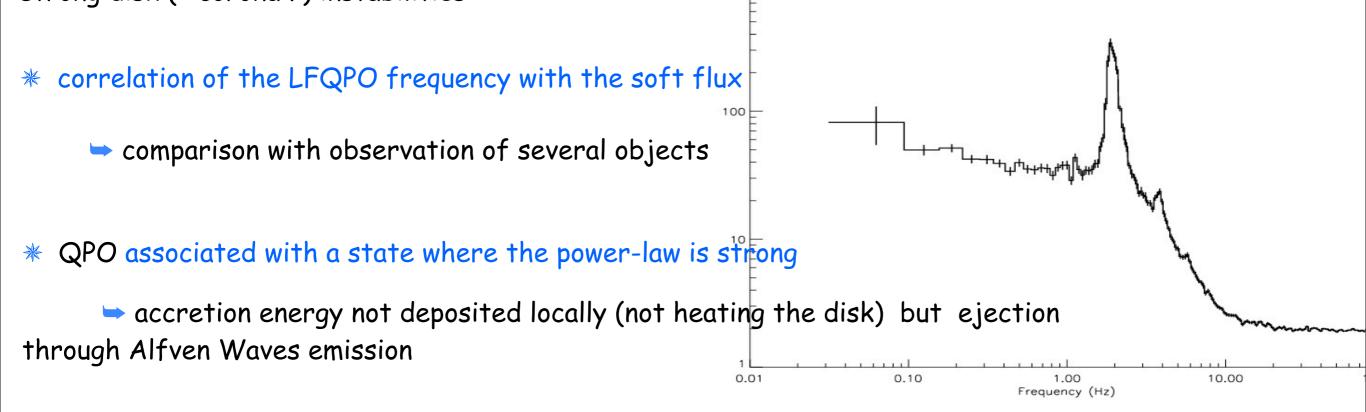
#### Looking for a LFQPO model



> the frequency of the spiral wave is a few time the rotation frequency at the inner edge of the disk

\* coherence  $\Rightarrow$  QPO can't be due to blobs  $\Rightarrow$  global organized motion of the gas = normal modes

\* rms amplitude as much as  $30\% \Rightarrow$  QPO can't happen by random processes (blobs)  $\Rightarrow$  must be due to strong disk (+ corona ?) instabilities



\* the existence of several "types" of LFQPO with distinct characteristics (A, B and C)

\* lag sometimes changing sign, sub-harmonic structure...

### The Magnetic Flood Scenario

the AEI is a candidate to explain the LFQPO and can also be linked with ejection.

**next step**: assume that  $AEI \Leftrightarrow QPO$  ... then try to understand ....

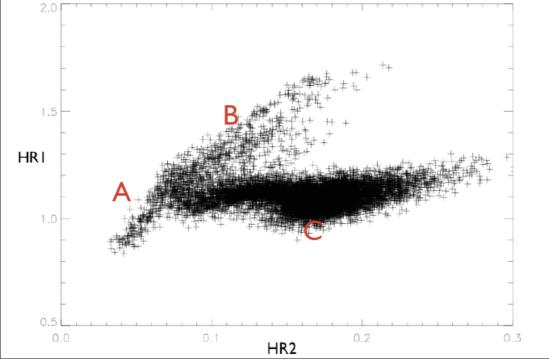
→ "Magnetic Floods" scenario : the cycles are determined by gradual accumulation and sudden destruction of magnetic flux

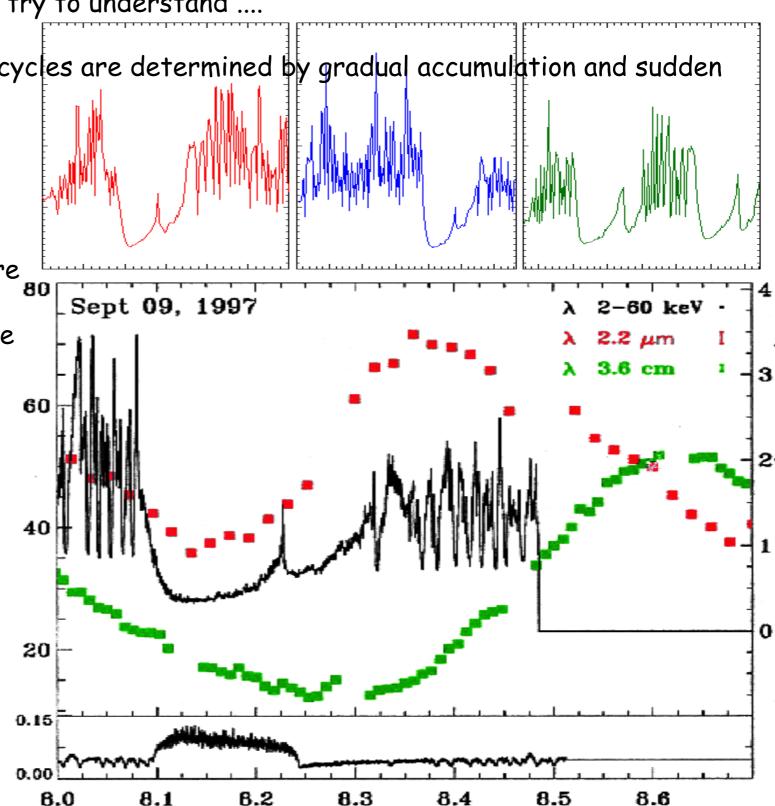
we looked at the cycles of GRS 1915+105 because:

- we had already worked on that source before 80

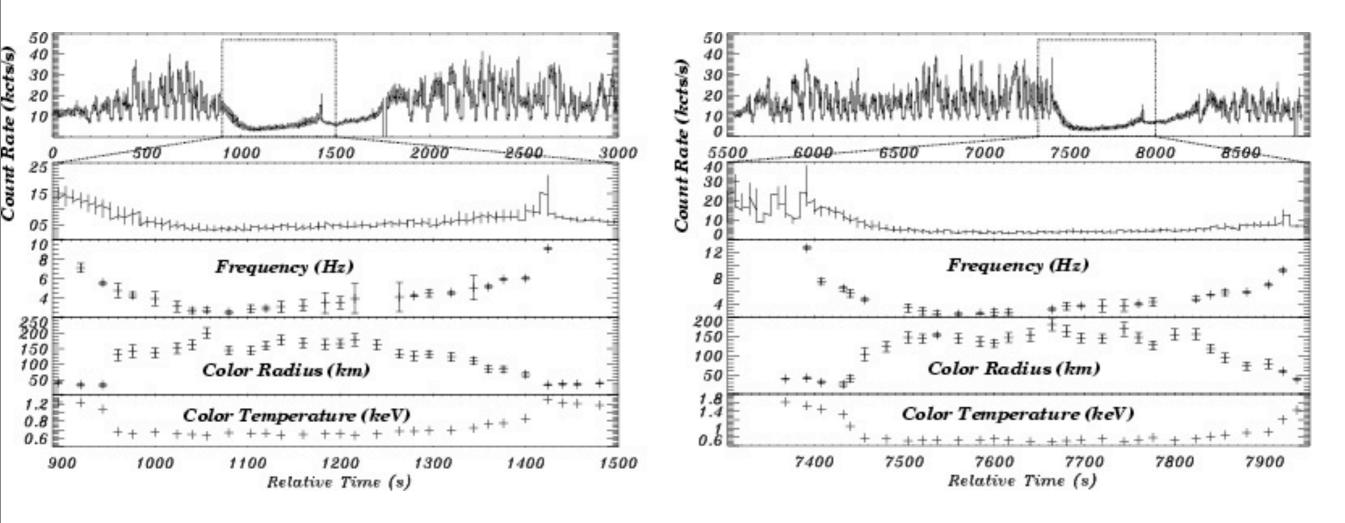
- multi-wavelength observations were available

- it has known (and highly repeatable) cycles between state C, A and B (Belloni et al 2000)





### The Magnetic Flood Scenario



The cycle begins in the "high/soft -thermal state" during which magnetic field is accumulated in a turbulent disk (possibly driven by the MRI)  $\rightarrow$  B becomes of the order of equipartition

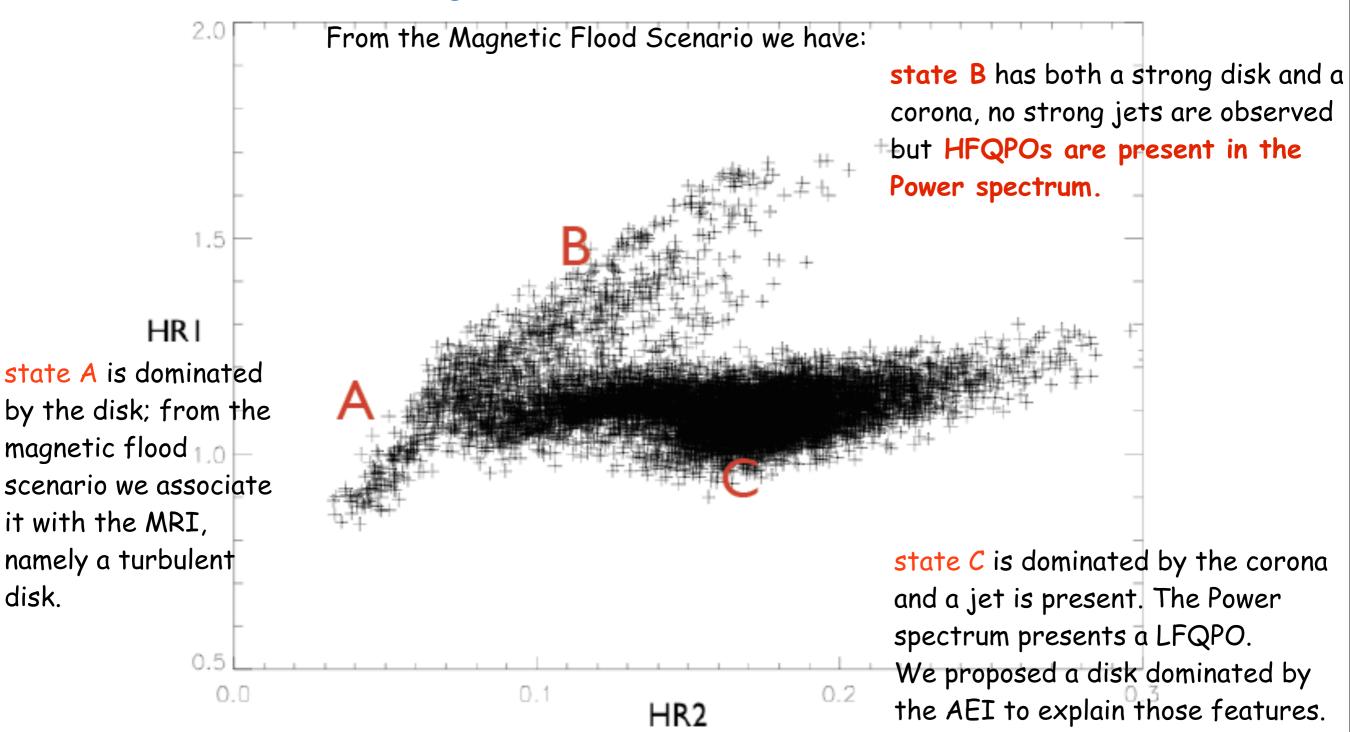
→ the disk becomes unstable to the AEI, LFQPO appears. it is in the "Low-Hard state"

 $\rightarrow$  at the intermediate peak we have reconnection of the magnetic field, causing an ejection and the plasma  $\beta$  to increase ending the AEI and the LFQPO causing the disk to go back to a softer state

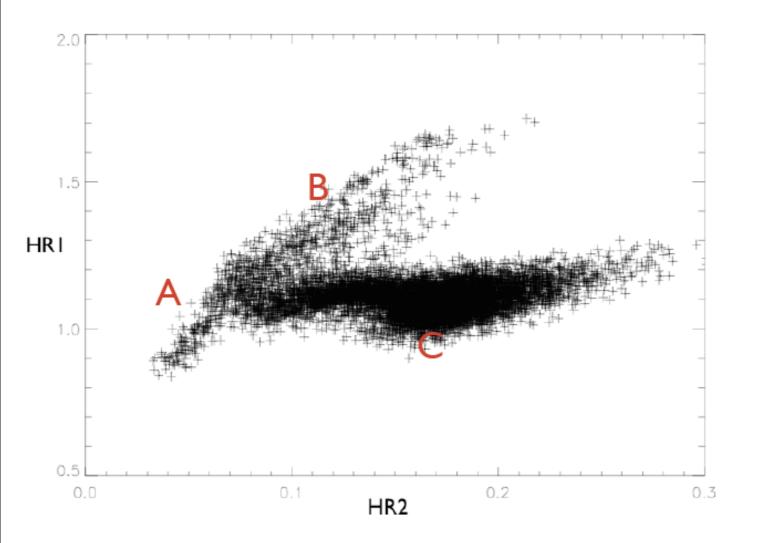
#### → we linked instabilities occurring in the disk with observational states

# What is at the origin of the three states?

All the observations from GRS 1915+105 can be classified in 12 classes made of three states labeled A, B and C as shown on the color-color diagram.



### What is at the origin of the three states?



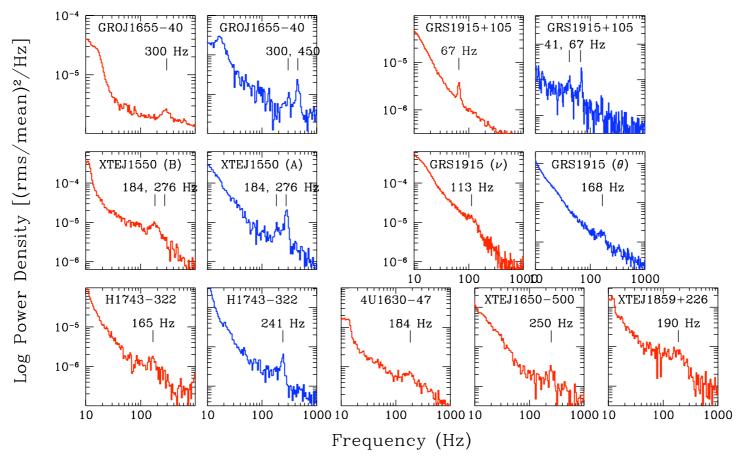
In 2006 we proposed the Rossby Wave Instability (RWI) to be at the origin of HFQPOs in black hole systems.

This instability requires an extremum of vortensity which exists in disks with their inner edge close to the last stable orbit.

One interesting characteristic of this instability is that the m=1 mode is not dominant but it is rather a mix of the m=2 and m=3 modes that dominates.

this instability was the last piece of the puzzle and we now have three disk instabilities which properties make them good candidates to explain the three states

#### Looking for a HFQPO model: what we need to explain



small variations in the observed frequency (maximun change in frequency of about +/- 15%)
 the existence of the RWI is linked to the position of maximun of vortensity

\* HFQPOs appear alone or in "pairs" (with related frequencies, most of the time in a 2:3 ratio, sometimes 1)

 $\Rightarrow$  need one mechanism that can select several linked frequencies depending on the disk conditions

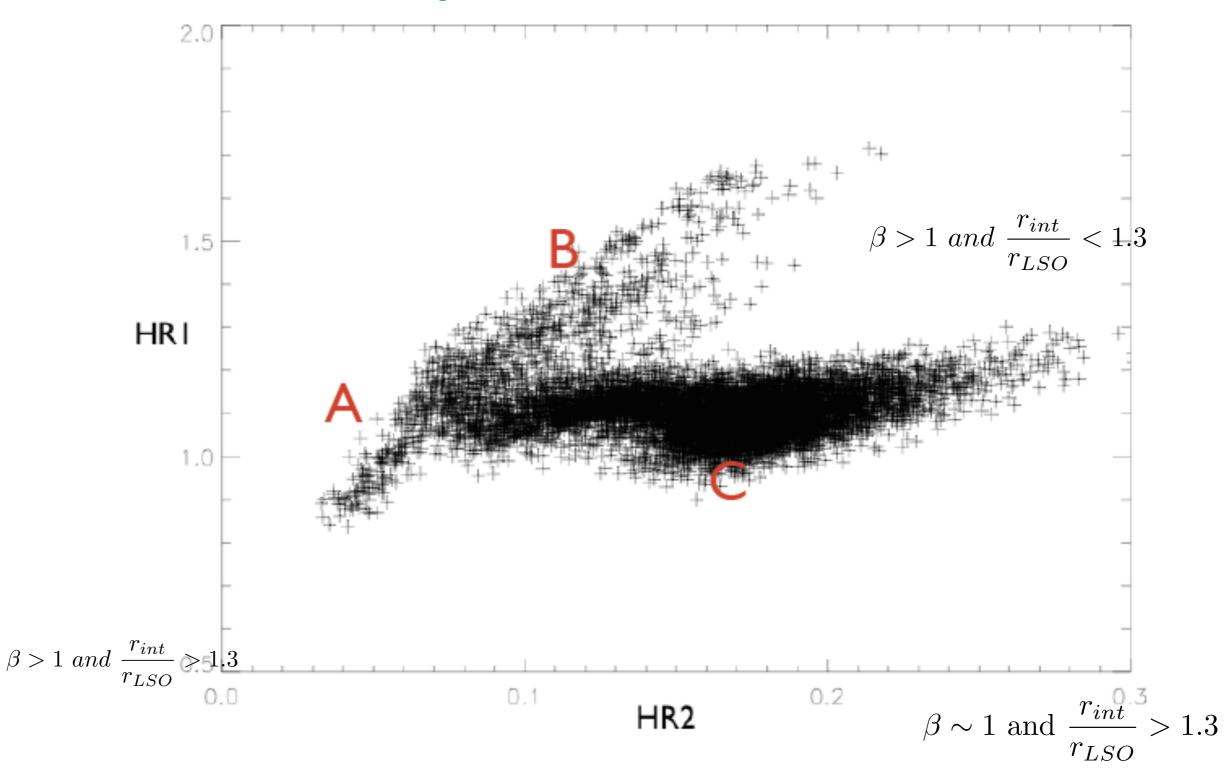
- ⇒ depending on the inner boundary, the dominant mode of the RWI is the m=2 / m=3, or the m=1
- \* HFQPOs can occur in the absence of LFQPOs
- \* when they co-exist we have "unusual" LFQPO (type A and B)

\* HFQPOs rms amplitudes are much lower than LFQPO and seem anti-correlated with the LFQPO rms

⇒ the HFQPO models need to be coherent with a LFQPO model as they have to co-exist in the disk while being independent

# The 3 instabilities at the origin of 3 states

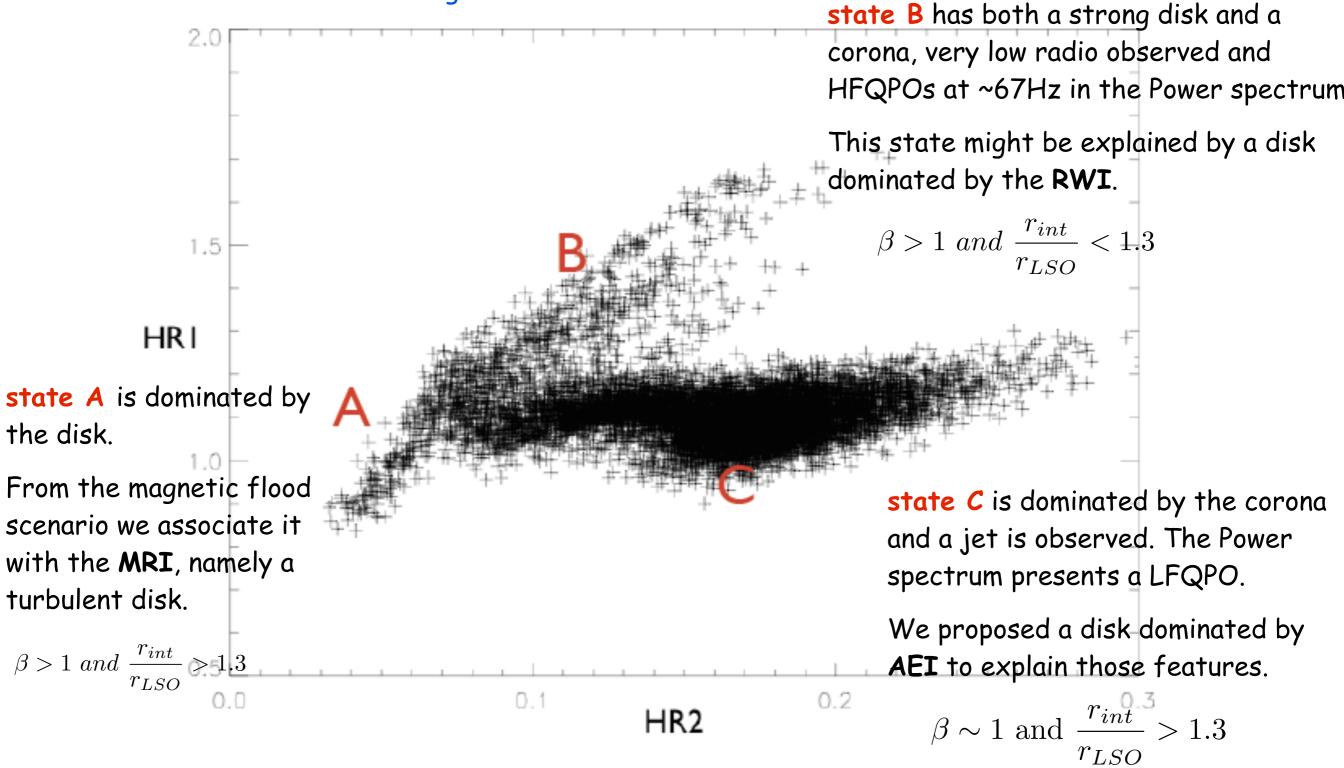
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# The 3 instabilities at the origin of 3 states

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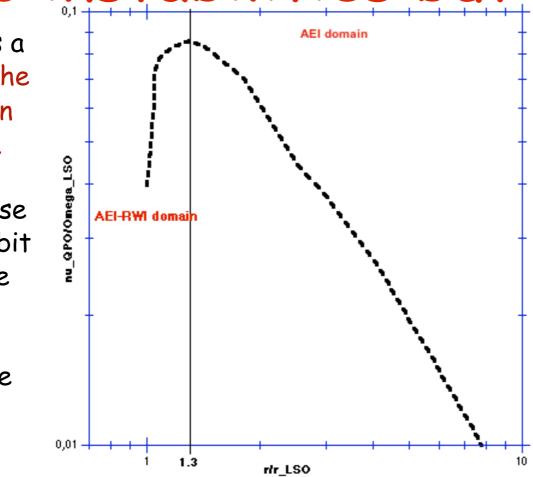
→ this can explain the observed 3 states of GRS 1915+105 but not the presence of both HFQPO and LFQPO

### 3 instabilities but 4 states

We see that there is a domain where both the AEI and the RWI can be active in the disk.

The inner edge is close to the last stable orbit and the disk is at the equipartition.

In that state we have both the LFQPO and the HFQPOs



validity

RWI: inner edge of the disk close to the last stable orbit and an extremun of:

 $\mathcal{L}_B = \frac{\kappa^2}{2\Omega} \frac{\Sigma}{B^2}$ 

AEI: equipartition between gas pressure and magnetic field and

$\partial$ ln	$\left(\kappa^{2}\Sigma\right)$	> 0
$\overline{\partial \ln r}$ III	$\left(\overline{2\Omega B^2}\right)$	> 0

we get the following four possible states with our three instabilities:

- MRI:	turbulent disk, similar to the A state, high/soft state, thermal state
- RWI:	hot disk, HFQPO, very low radio, similar to B state, other?
- AEI:	cold disk, LFQPO, jet, BLN, similar to the C state, low-hard state
- AEI-RWI:	warmer/hot disk, low radio, HF + LFQPOs, steep power-law, VHS

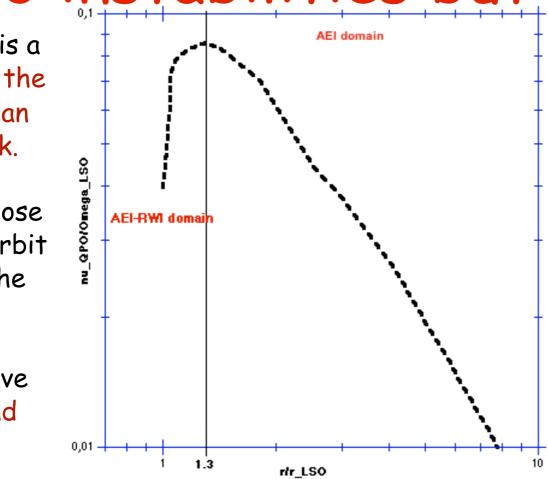
→ The fourth state exhibits both LF and HF QPOs, which occurs frequently in sources like XTE J1550-564.... It has since been observed GRS 1915+105 has become the first object to show all 4 states. The RWI alone state has not been observed yet in other microquasars.

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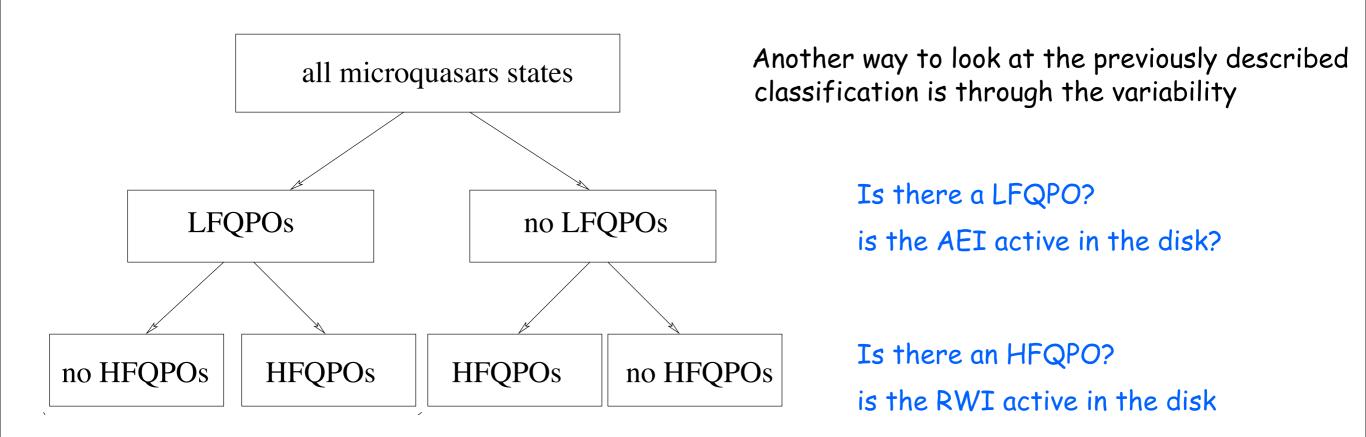
$$\frac{\partial}{\partial \ln r} \ln \left( \frac{\kappa^2 \Sigma}{2\Omega B^2} \right) > 0$$

we get the following four possible states with our three instabilities:

- MRI: $\beta > 1$  and  $\frac{r_{int}}{r_{LSO}} > 1.3$  turbulent disk, similar to the A state, high/soft state, thermal state...- RWI: $\beta > 1$  and  $\frac{r_{int}}{r_{LSO}} < 1.3$  hot disk, HFQPO, very low radio, similar to B state, other?- AEI: $\beta < 1$  and  $\frac{r_{int}}{r_{LSO}} > 1.3$  cold disk, LFQPO, jet, BLN, similar to the C state, low-hard state- AEI: $\beta < 1$  and  $\frac{r_{int}}{r_{LSO}} > 1.3$  cold disk, LFQPO, jet, BLN, similar to the C state, low-hard state- AEI-RWI: $\beta < 1$  and  $\frac{r_{int}}{r_{LSO}} < 1.3$  warmer/hot disk, low radio, HF + LFQPOs, steep power-law, VHS- The fourth state exhibits both LF and HF QPOs, which occurs frequently in sources like XTE

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## Variability based classification



This way to look at the classification is model-independent. It is not important what instability is at the origin of the LFQPO and the HFQPO as long as they can operate simultaneously in the disk as required by observation

#### another requirement for both LF and HF QPOs model

An interesting by-product of this classification is that it links together the behavior of the LFQPO with the HFQPO and could explain the LFQPO types in microquasars like XTE J1550-564.

## A quick look at LFQPOs type

When we are in a state where both the RWI and the AEI exist, it is actually the relativistic AEI which is present. The R-AEI has slightly different properties.

We did numerical simulations of a disk having both the AEI and the RWI and it appears that the dominant mode is always the same for both of the instabilities.

In that picture we have

C-type is the "regular AEI" A and B-type is the R-AEI with higher modes dominating

link between the type of LFQPO and the HFQPO mode

both instabilities (AEI and RWI) are in the same disk, the same modes are dominant for both instabilities.

(sub)-harmonic content (A- and B-types sometimes present sub-harmonics or "almost harmonics")

→ contrary to the non relativistic case, it is not the m=1 mode of the AEI which always dominates when the disk gets close to the last stable orbit. The non-dominating m=1 could be taken for a sub-harmonic. Also, the different modes of the AEI are in close-harmonic relationship, not exact.

the frequency of the A- and B-type seems to have a small range (compared to the C-type)

to be in the AEI-RWI state, the inner edge of the disk needs to be close to the last stable orbit so the LFQPO cannot change frequency as much as in the C-type case

### Conclusion

- using a consistent model : control by the gradual accretion of vertical magnetic flux, stored in the BH magnetosphere and destroyed in reconnection events (= ejections) when disk magnetic flux and stored flux are opposite
- we proposed a new classification based on which instability dominates the disk: equivalent to a variability based classification
- studying the instability behaviour in that new state gives a possible explanation for LFQPO types which needs to be further tested
- for the next step is to look at the consequences of the different instabilities on ejection to also link the four states with the radio emission observed.
- ø we are also working on making a list of observational constraints against which to test both LFQPO and HFQPO models

# Alfven Wave Emission Toward the Corona

The Rossby vortex twists the footpoint of the field lines threading the disk. If the disk has a low density corona:

> emission of Alfven Waves twisting

• energy and angular momentum extracted from the disk will be transferred to the corona where they can power a wind or a jet

We describe this via a variational form: ->

F = energy of the waves + i ( outgoing spiral

+ coupling with the vortex + kz Alfven Waves)

imaginary term  $\Leftrightarrow$  amplification or damping of the wave

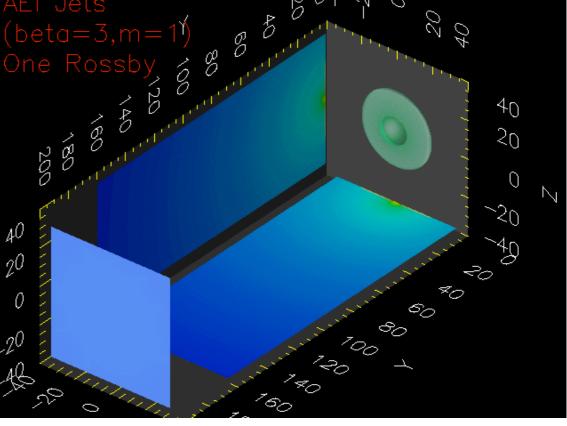
the Alfven terms are singular at the vortex radius

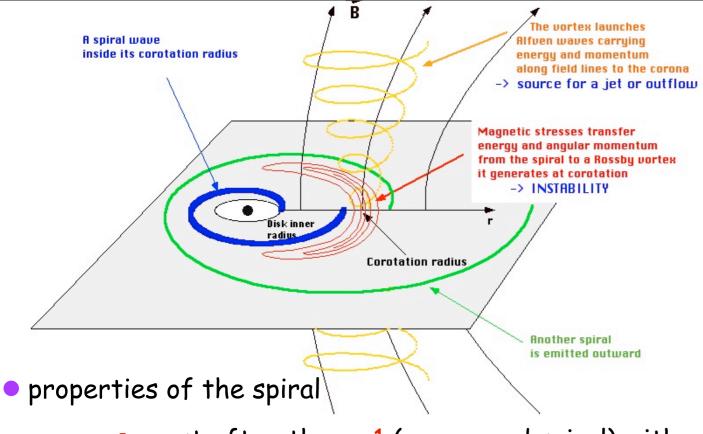
Typical Xray Binary aspect ratio  $h/r \sim 10^{-2}$ 

Efficiency of the mecanism: we compute the ratio of the flux emitted toward the corona as Alfven waves to the energy removed from the inner region (flux at the corotation radius)

 $\frac{F_{Alfven}}{E} \sim \left(\frac{\rho_{corona}}{\rho_{corona}}\right)^{1/2} \left(\frac{r}{h}\right)^{3/2}$ 





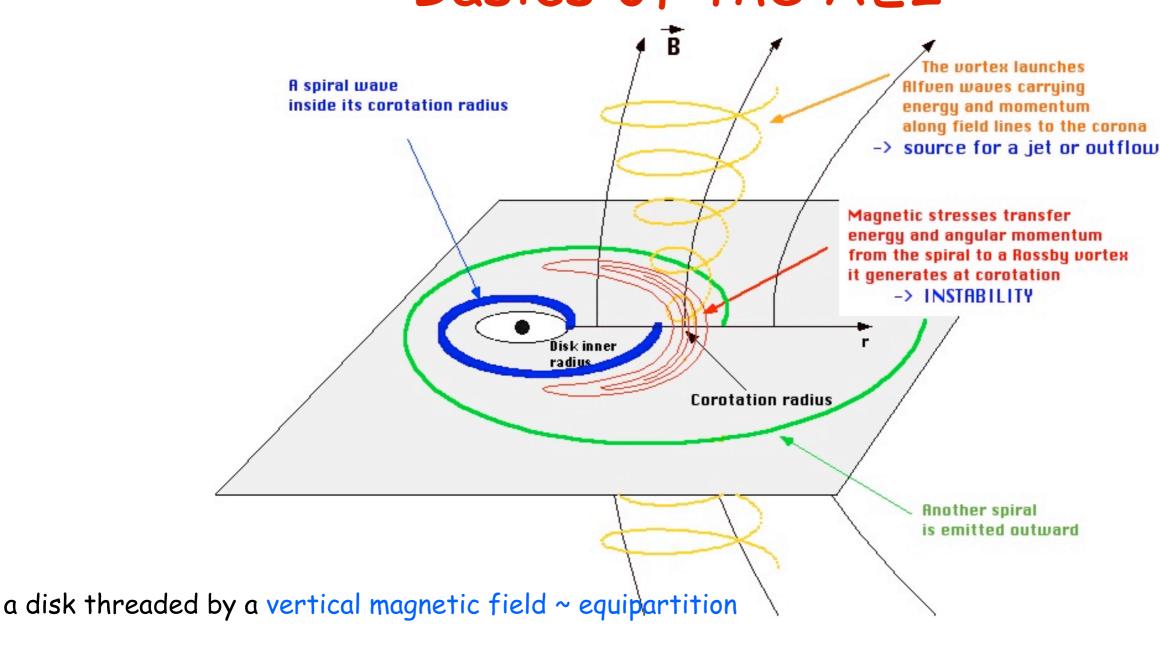


# Basics of the AEI

→ most often the m=1 (one armed spiral) with  $\omega \sim 0.1 - 0.3 \Omega_{int}$  (rotation frequency at the inner edge of the disk)

- differential rotation + differential vorticity
  - Instable by coupling to a Rossby vortex (~ great red spot of Jupiter)
- extracts energy and angular momentum from the disk (> accretion) and stores them in the Rossby vortex
- if the disk has a low density corona

### Basics of the AEI



→ large scale instability

a spiral wave ~ galactic spirals but driven by magnetic stresses rather than self-gravity

same structure of large scale normal modes (= standing patterns) as in galaxies

The perturbation is almost constant across the disk height

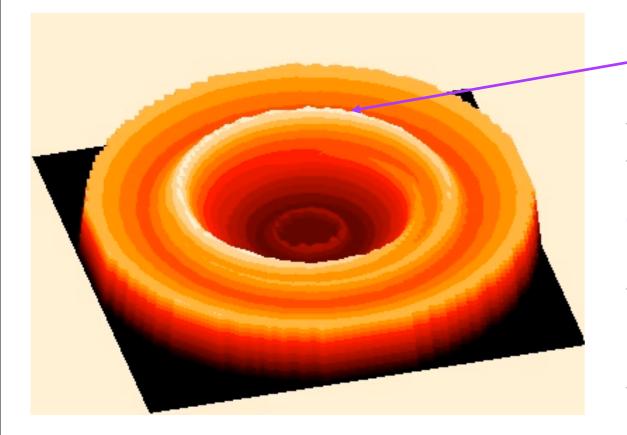
→ 2D simulation possible

All self similar MHD jet models (Blandford & Payne, Pelletier & Pudritz, ...) are unstable

### a modulation of the disk emission

work done with M.Muno, L.Prat

The effects of the spiral shock is the creation of a hot-point in the disk



 $\rightarrow$  and also a thickenning of the disk

From 2D non-linear MHD simulation we get the following thickness (with a simplistic model for disk height and temperature)

Computing the X-ray emission from such an accretion disk shows a modulation in the flux of a few percent (5 to 10%)

→ we need to improve the computation of the disk thickness to obtain a better estimate of the rms

what does it gives:

- geometric modulation of the flux
- hot region in a warn disk  $\rightarrow$  growth of the rms amplitude with energy
- not only the disk but the jet will obscur the hot-point

 $\rightarrow$  a way to explain the reverse time lag

### The AEI as a model for the LFQPO

#### \* frequency between 1-10 Hertz

Frequency of the spiral wave  $\omega \sim 0.1 - 0.3 \Omega$  int (rotation frequency at the inner edge of the disk)

#### \* stability in time

Iarge scale structure as in galaxies (quasi-standing spiral structure)

#### \* rms amplitude as much as 30%

we observed in the simulation as much as 10% in a oversimplified model, working on a better approximation

#### \* correlation with the soft flux

comparison with observation in Varniere et al A&A 2002 and Mikles et al. ApJ 2009

#### \* QPO associated with a state where the power-law always dominates

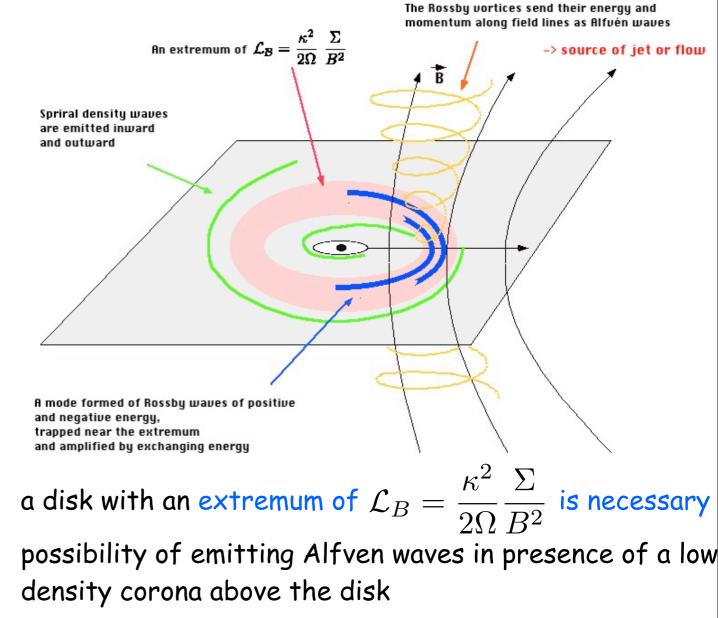
accretion energy not deposited locally (not heating the disk) but ejection through Alfven Waves emission

\* lag sometimes changing sign, sub-harmonic structure

possibility of geometrical effects from the jet

### RWI at the origin of the state B?





In 2006 we proposed the Rossby Wave Instability (RWI) to be at the origin of HFQPOs in black hole systems.

HR2

0.2

0.3

0.1

1.5

1.0

0.0

HRI

This instability requires having an extremum of vortensity which exists in disks with their inner edge close to the last stable orbit.

One interesting characteristic of this instability is that, depending on the inner boundary, the m=1 mode is not dominant but it is rather a mix of the m=2 and m=3 modes that dominates.

→ we need to see if that instability is a good candidate for the HFQPOs in microquasars and therefore a candidate to explain the state B of GRS 1915+105

### The AEI as a model for the LFQPO

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Frequency of the spiral wave  $\omega \sim 0.1 - 0.3 \Omega$  int (rotation frequency at the inner edge of the disk)

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### RWI: toward a model for the HFQPO?

#### \* small variation in the observed frequency

+ the existence of the RWI is linked to the position of maximum of vortensity

#### \* seems to select pairs of frequencies, in a 3:2 ratio

⇒ depending on the inner boundary, the dominant mode of the RWI is the m=2 / m=3, or the m=1

#### \* the HFQPO model need to be coherent with a LFQPO model

+ the condition for the magnetized RWI are similar to the one of the relativistic AEI

#### \* link with the appearance of unusual (type A and B) LFQPO

➡ We proposed an explanation for the different types of LFQPO based on the behavior of the relativistic AEI

\* we still have several points to work on before being able to compare the RWI-HFQPO model to observations, especially the computation of the observed rms amplitude, the Alfven wave emission, the variation of frequency and 3D MHD simulation...

 $\Rightarrow$  We are also working on better understanding the constraints from observations.

+ this instability was the last piece of the puzzle and we now have three disk instabilities which could explain the three states

