

Testing the No-Hair Theorem with Quasi-Periodic Variability

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The No-Hair Theorem

The Kerr solution is the only possible stationary axisymmetric metric in vacuum that has no naked singularities and no closed time-like loops.

- Black holes have only 2 parameters: mass and spin
- Any other signature radiated away by gravitational waves
- Expectation: All astrophysical black holes are Kerr black holes

However...

- Non-spherical collapse: naked singularities possible
Shapiro & Teukolsky 1992, PRD, 45, 44
- Kerr solution not unique to GR
Psaltis et al. 2008, PRL, 100, 091101

Testing the No-Hair Theorem Observationally



Kerr black hole?

Something else?

- Expand spacetime in multipole moments (mass M , spin a , ...)
- Parameterize quadrupole moment: $q = -(a^2 + \varepsilon)$
- Use observations to measure (at least) 3 moments
- Check whether: $q = -a^2$ i.e.: $\varepsilon = 0$

Ryan 1995, 1997a,b
Barack & Cutler 2004, 2007
Collins & Hughes 2004
Glampedakis & Babak 2006
Brink 2008, 2009
Gair et al. 2008

Li & Lovelace 2008
Will 2008
Vigeland & Hughes 2010
Johannsen & Psaltis 2010a, 2011b
Vigeland et al. 2011

A New Metric for Rapidly Spinning Black Holes

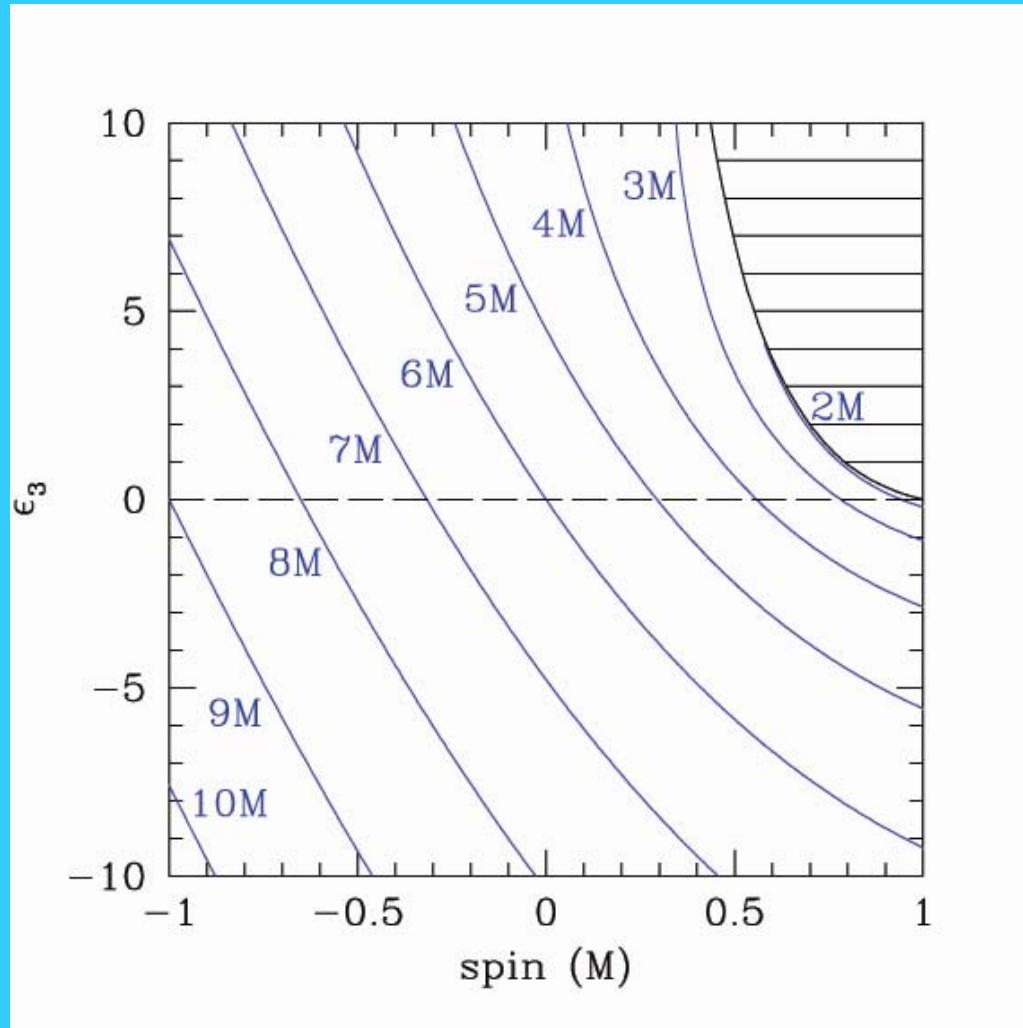
$$\begin{aligned} ds^2 = & -\left(1 - \frac{2Mr}{\Sigma}\right)(1+h(r,\theta))dt^2 - \frac{4aMr\sin\theta}{\Sigma}(1+h(r,\theta))dtd\phi \\ & + \frac{\Sigma(1+h(r,\theta))}{\Delta + a^2h(r,\theta)\sin^2\theta}dr^2 + \Sigma d\theta^2 \\ & + \left[\sin^2\theta\left(r^2 + a^2 + \frac{2a^2Mr\sin^2\theta}{\Sigma}\right) + h(r,\theta)\frac{a^2(\Sigma + 2Mr)\sin^4\theta}{\Sigma}\right]d\phi^2 \end{aligned}$$

where

$$h(r,\theta) \equiv \mathcal{E}_3 \frac{M^3}{r^3}$$

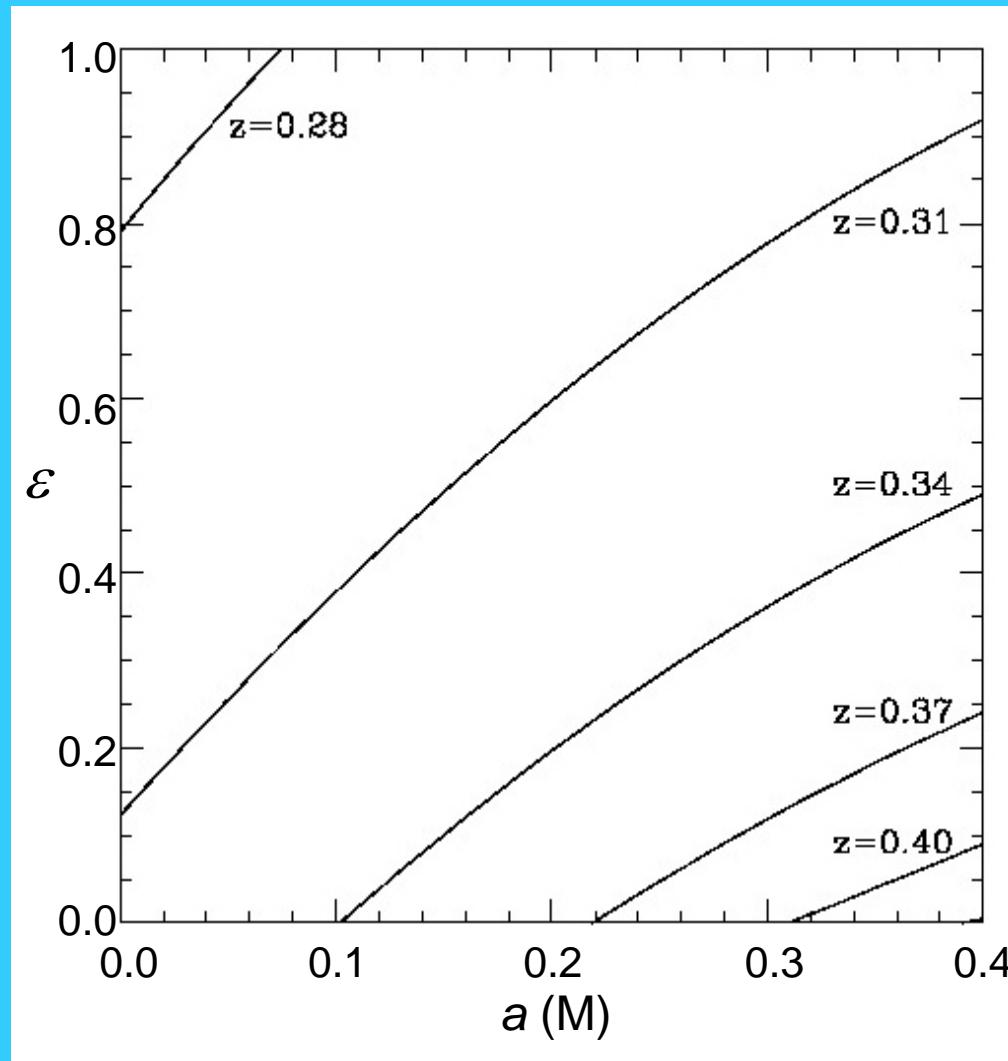
Johannsen & Psaltis 2011b, PRD, 83, 124015

Quadrupole Effects: (i) location of the ISCO



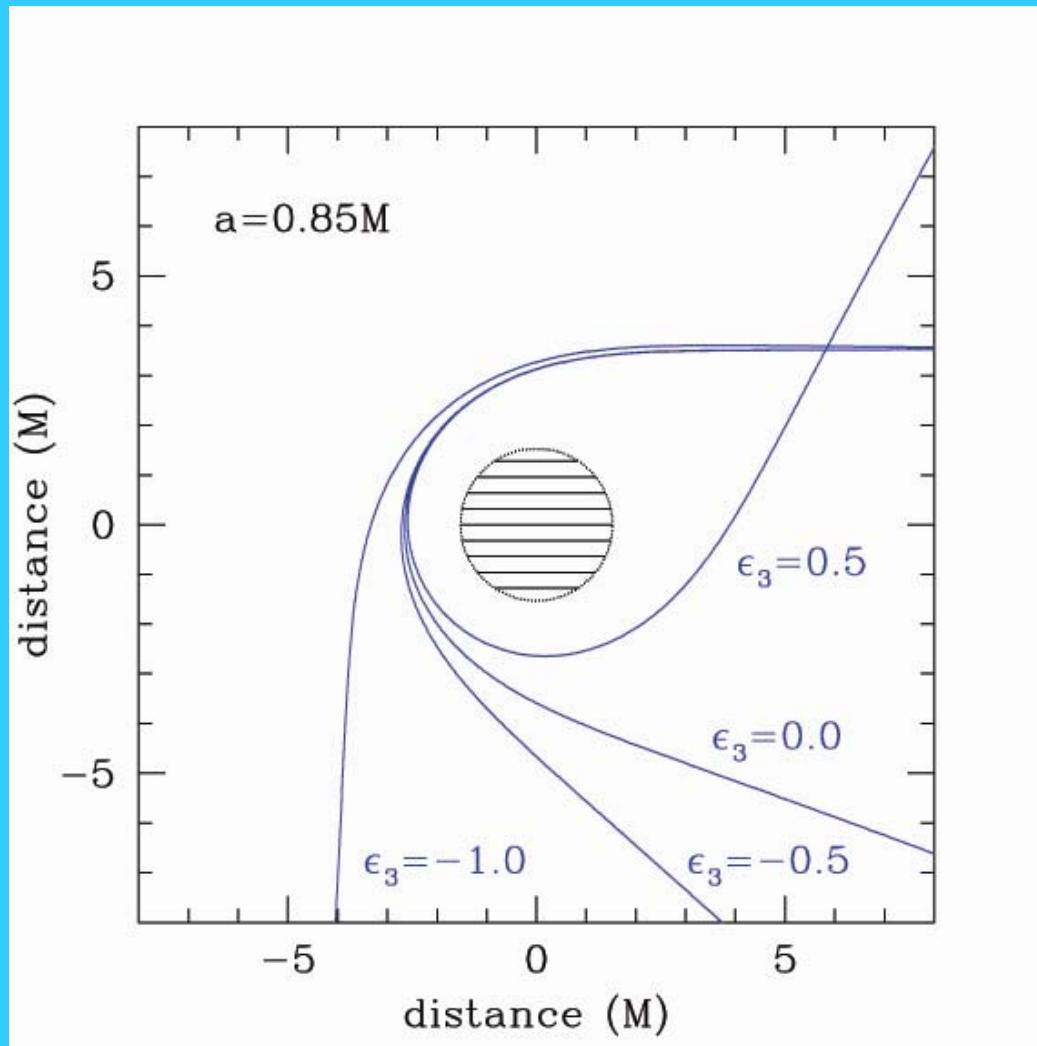
Johannsen & Psaltis 2011b, PRD, 83, 124015

Quadrupole Effects: (ii) photon redshift



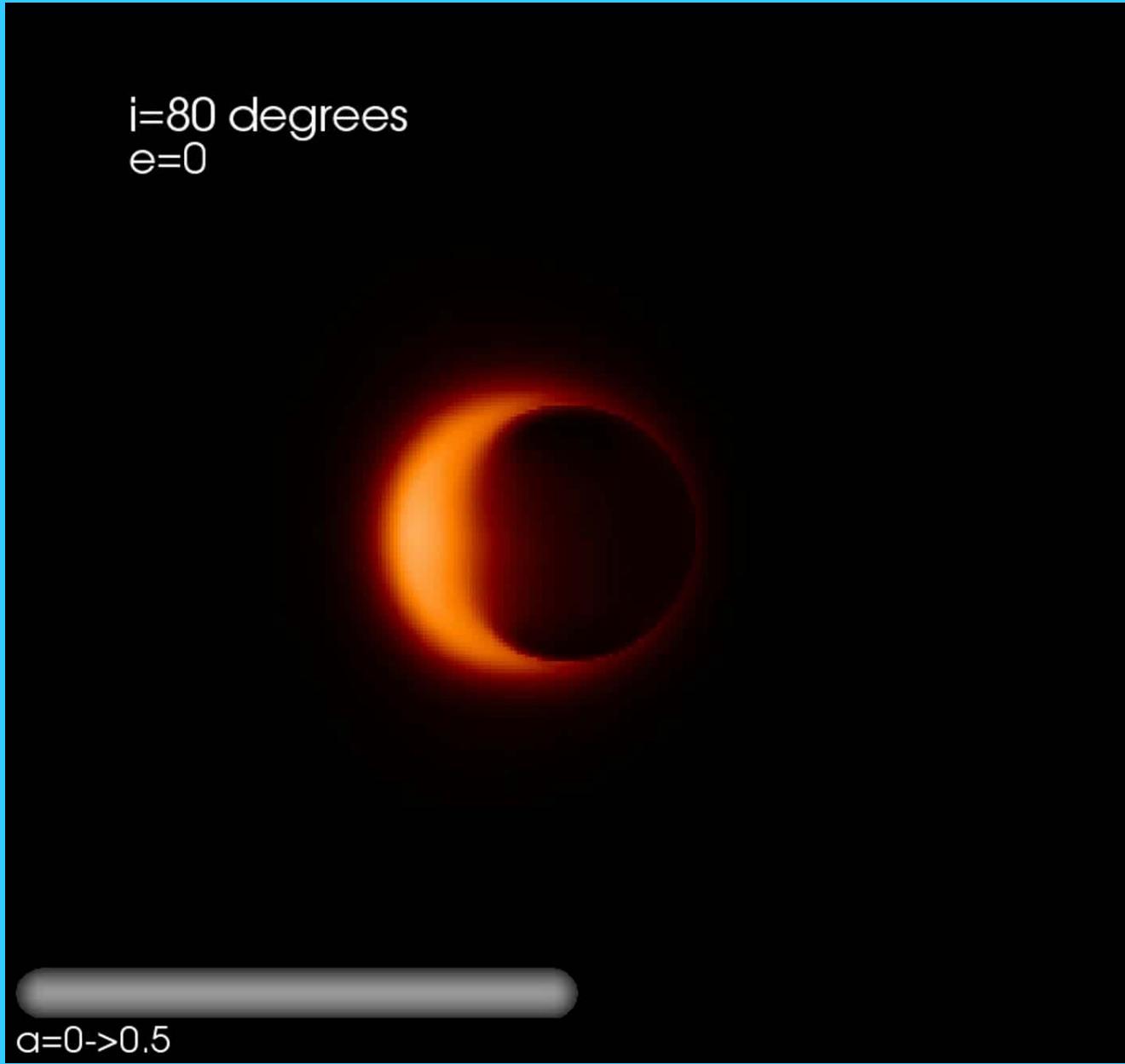
Johannsen & Psaltis 2010a, ApJ, 716, 187

Quadrupole Effects: (iii) gravitational lensing

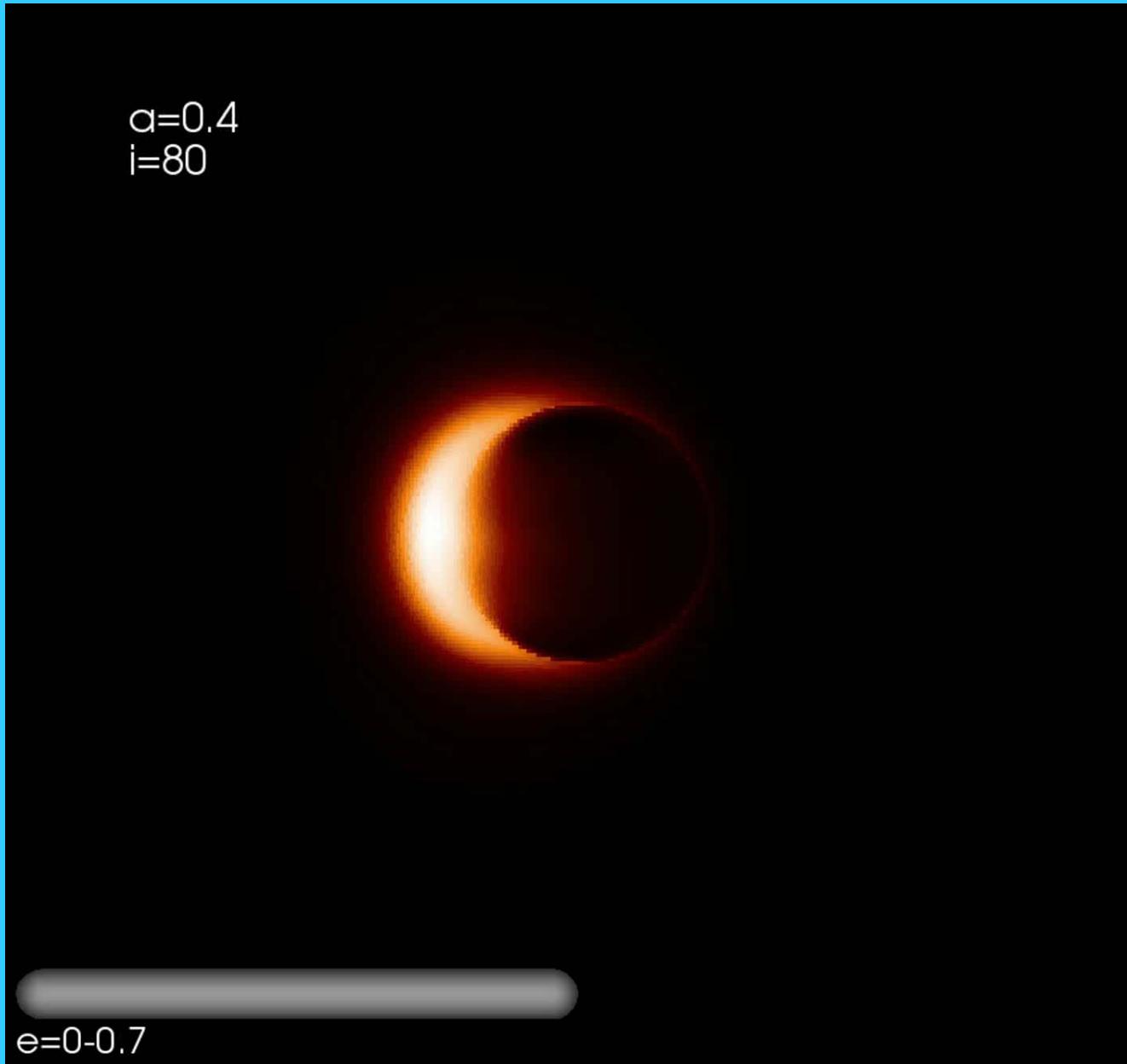


Johannsen, arXiv:1105.5645

I. Imaging the Shadow of Sgr A*

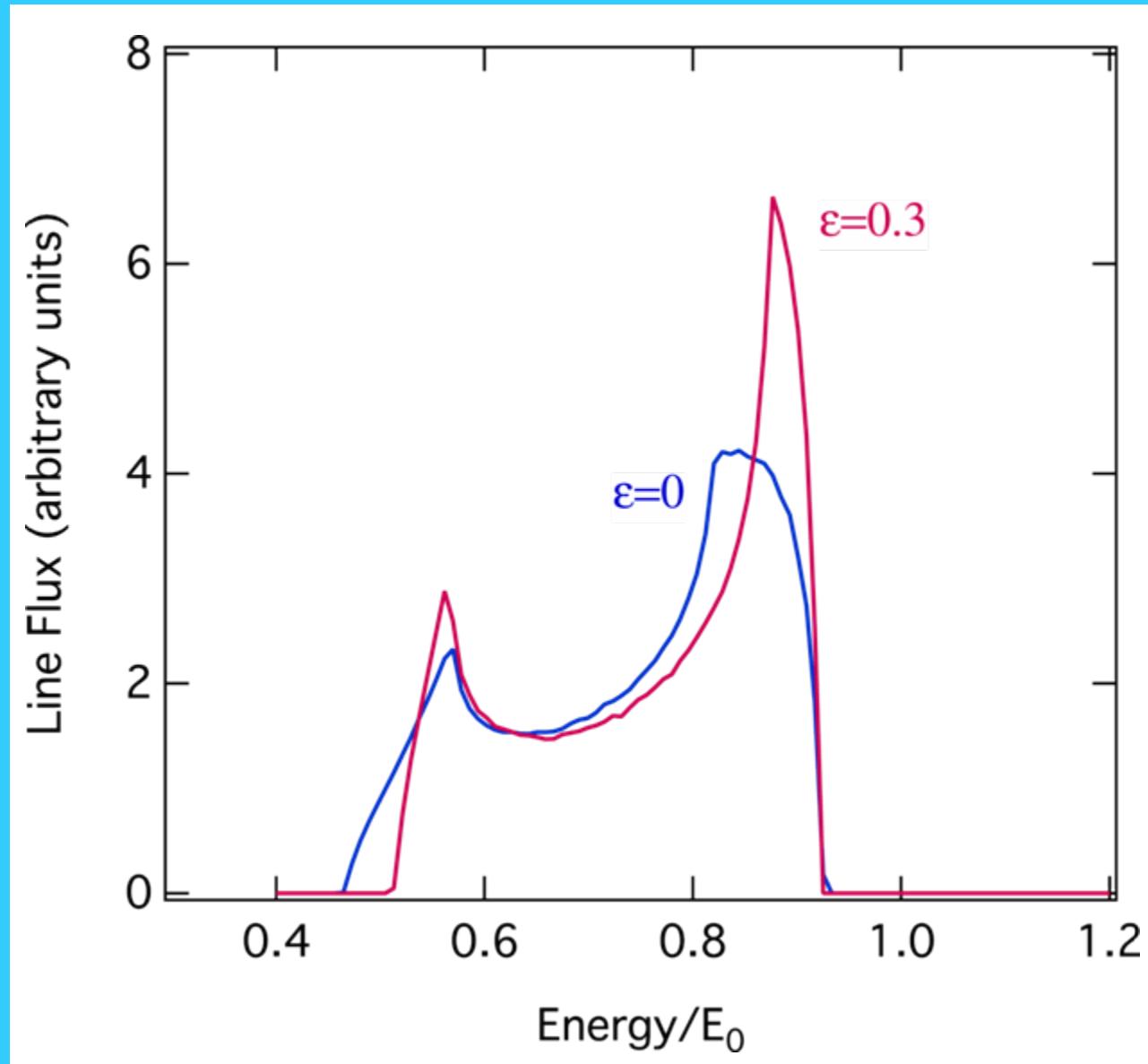


with A. Broderick, A. Loeb
Johannsen & Psaltis 2010a, ApJ, 716, 187



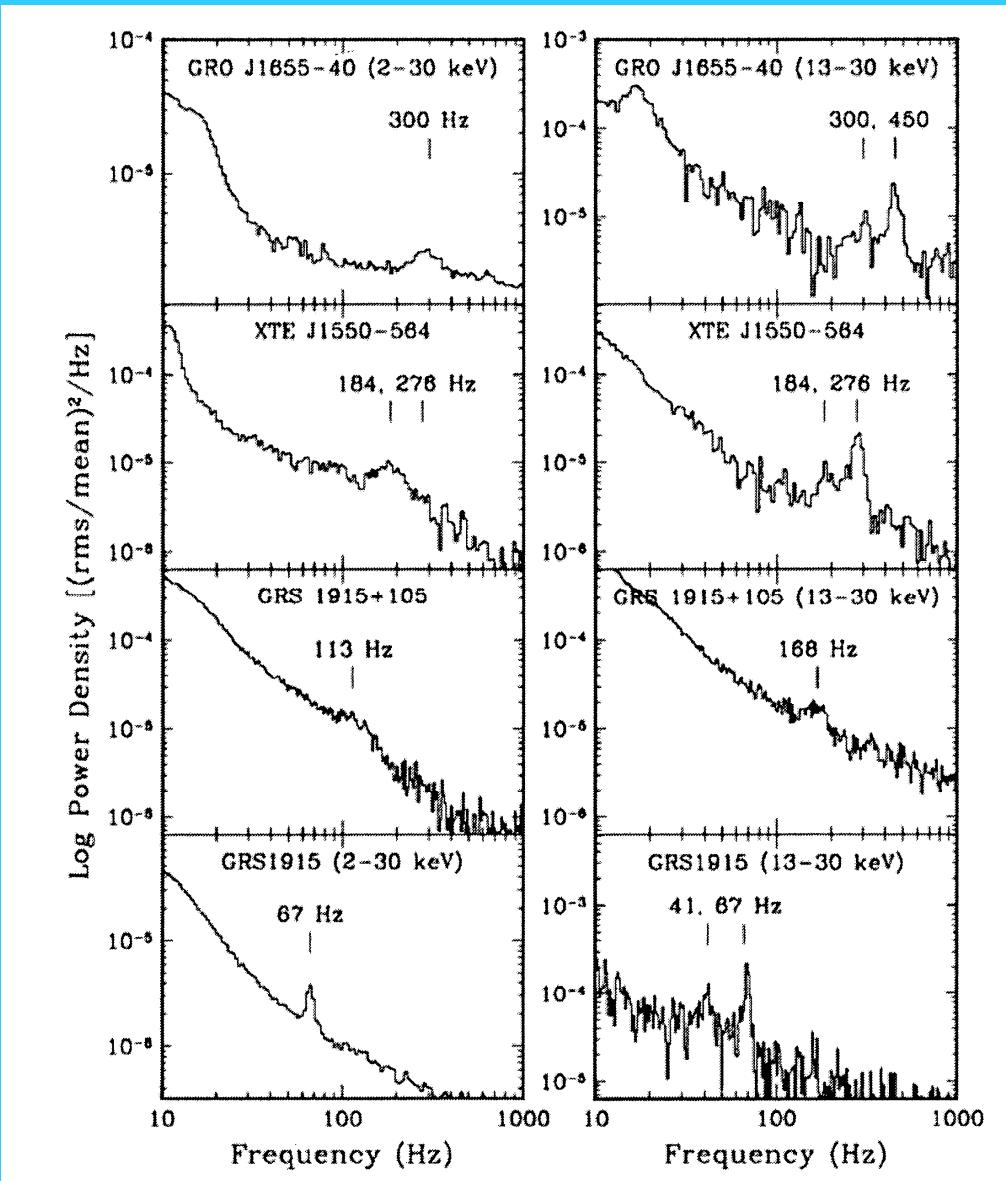
with A. Broderick, A. Loeb
Johannsen & Psaltis 2010a, ApJ, 716, 187

II. Fluorescent Iron Lines



Psaltis & Johannsen, arXiv:1011.4078

III. Quasi-Periodic Variability



McClintock & Remillard 2006, in *Compact Stellar X-Ray Sources*

1. Diskoseismology Model

Perez et al. 1997, ApJ, 476, 589
Silbergleit et al. 2001, ApJ, 548, 335

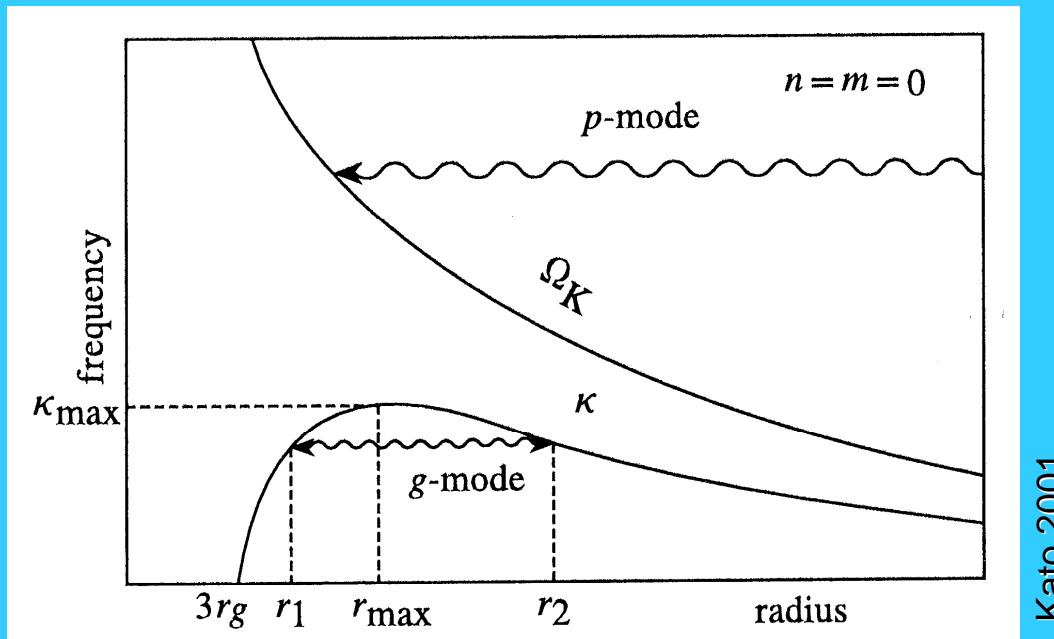
Consider perturbations on thin accretion disk:

Continuity eq.
Energy eq.
Eqs. of motion

} dispersion relation



g-modes
c-modes

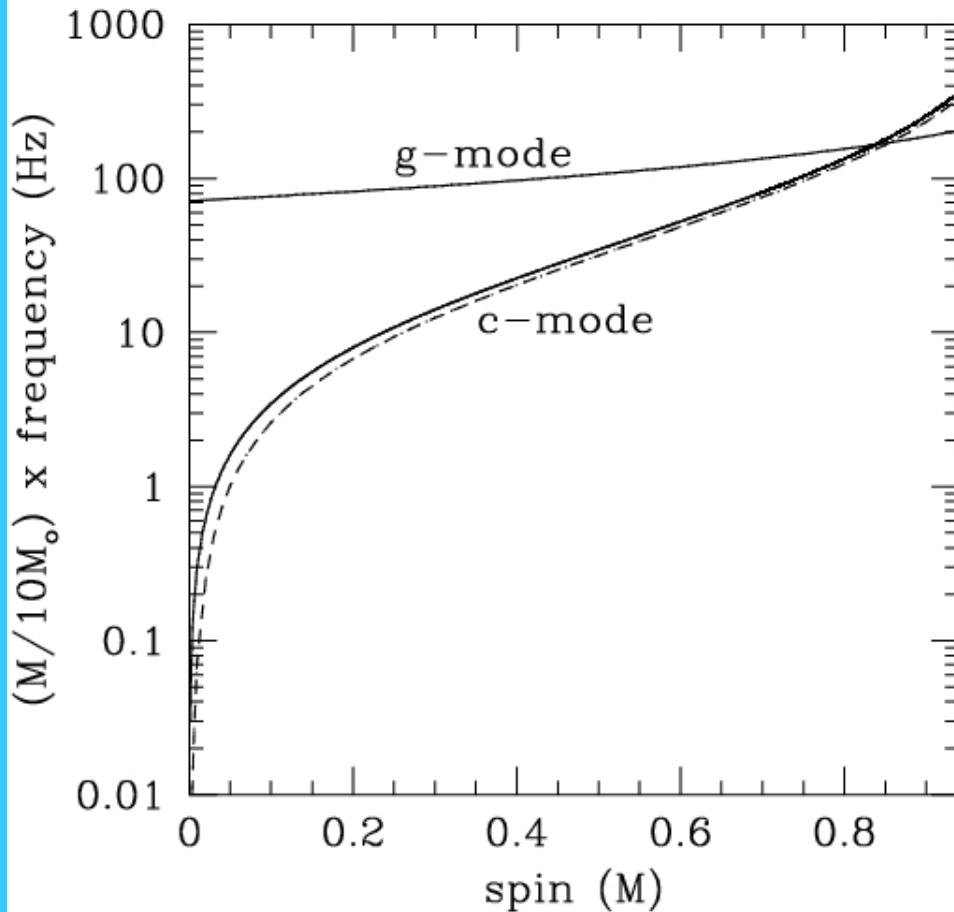


fundamental g-mode:

$$\frac{\kappa_{\max}}{2\pi}$$

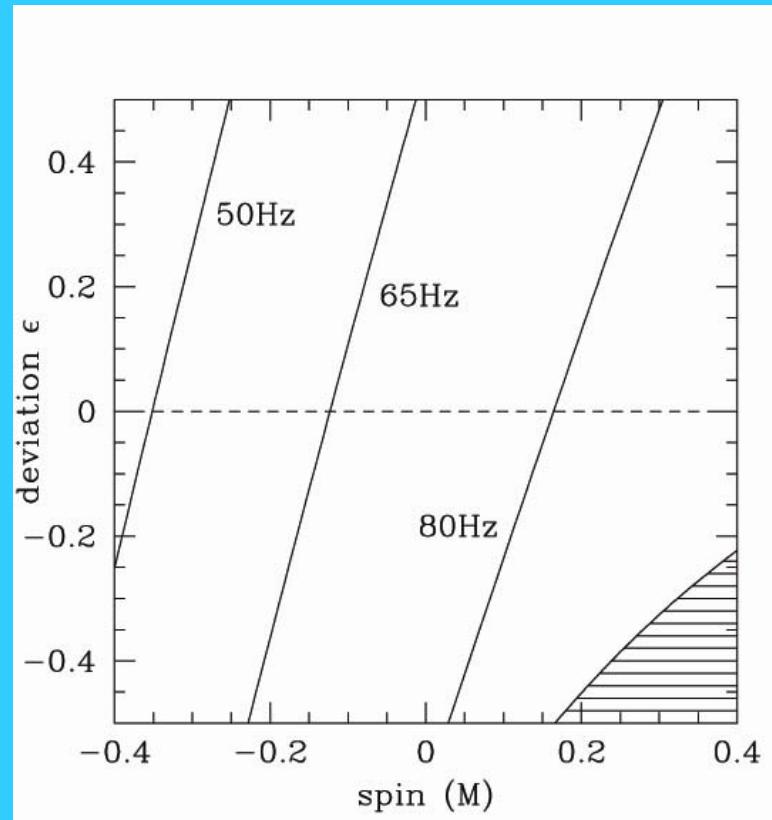
fundamental c-mode:

$$\frac{\Omega_{LT}(\text{ISCO})}{2\pi}$$

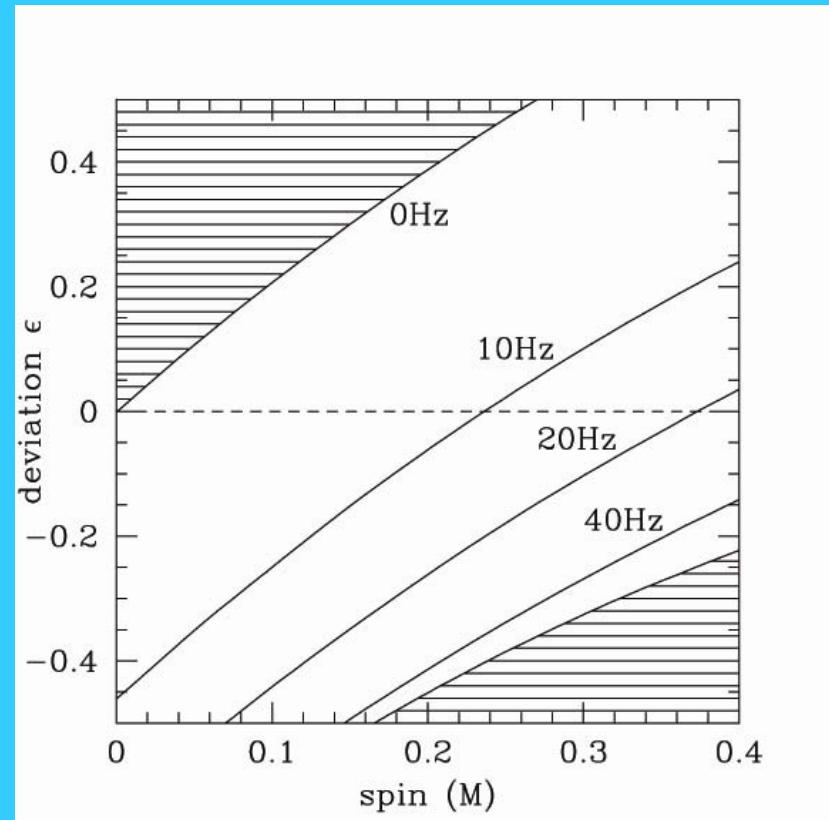


Kerr black hole

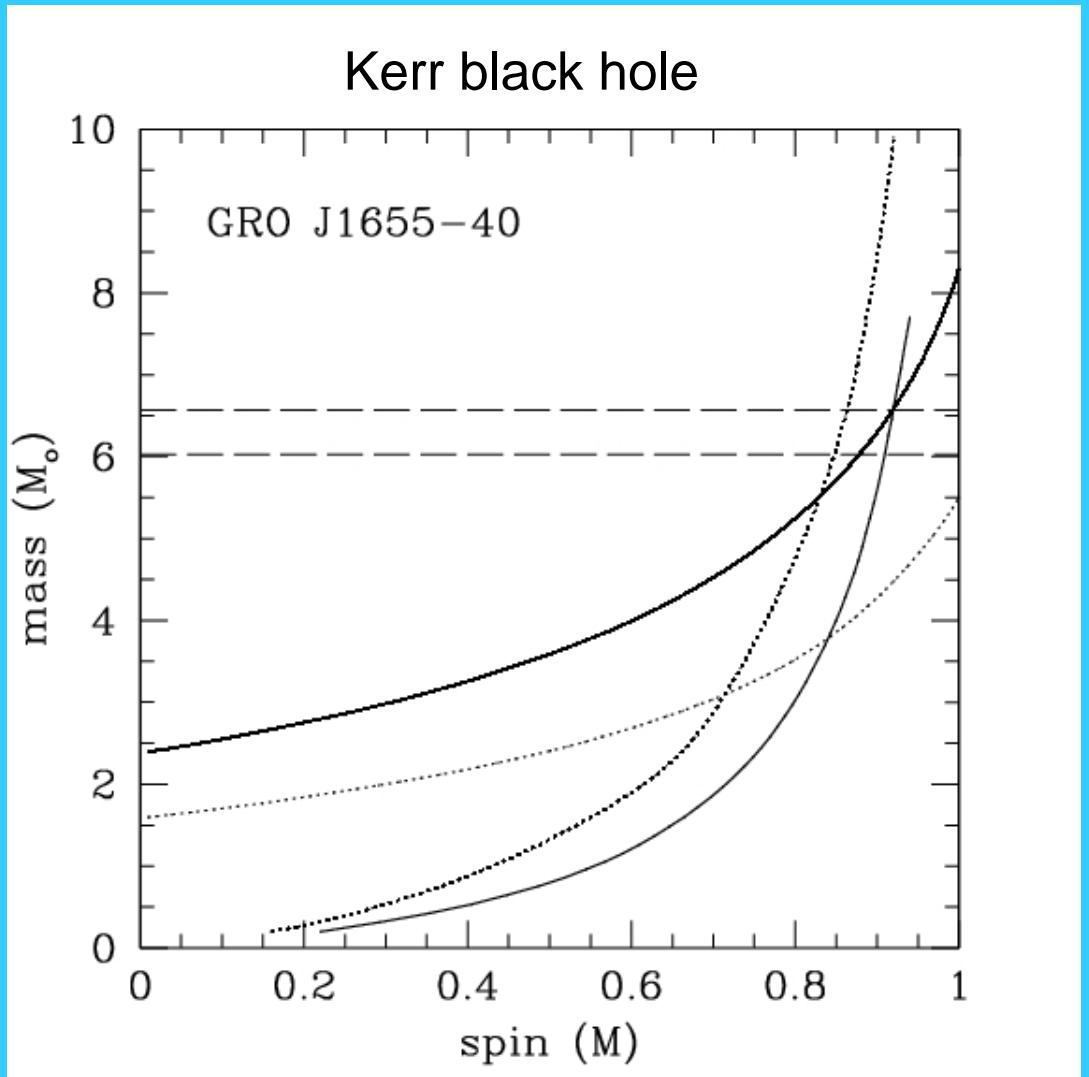
Wagoner et al. 2001, ApJ, 559, L25
Johannsen & Psaltis 2011a, ApJ, 726, 11



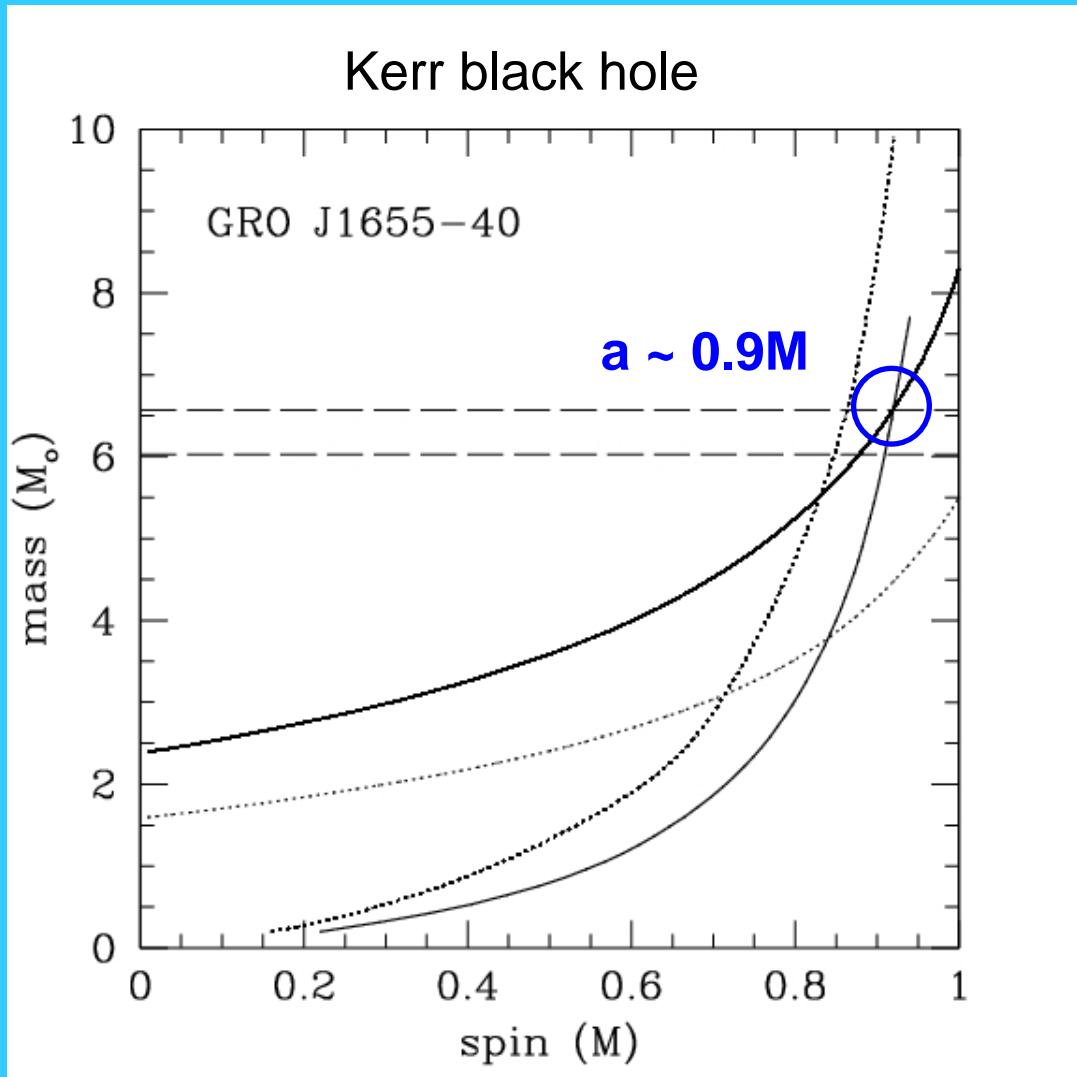
fundamental g-mode
(in units of 10 solar mass BH)



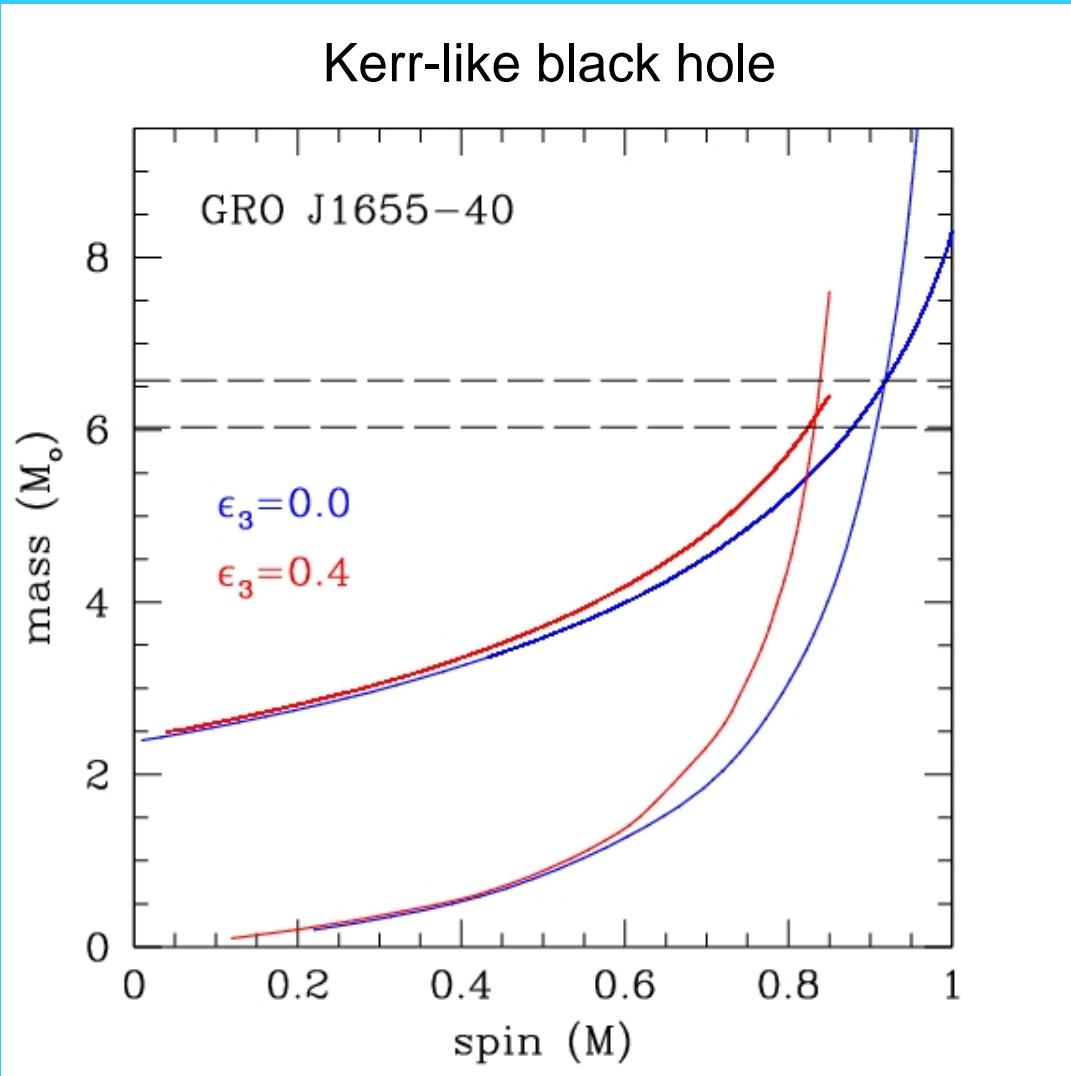
fundamental c-mode
Johannsen & Psaltis 2011a, ApJ, 726, 11

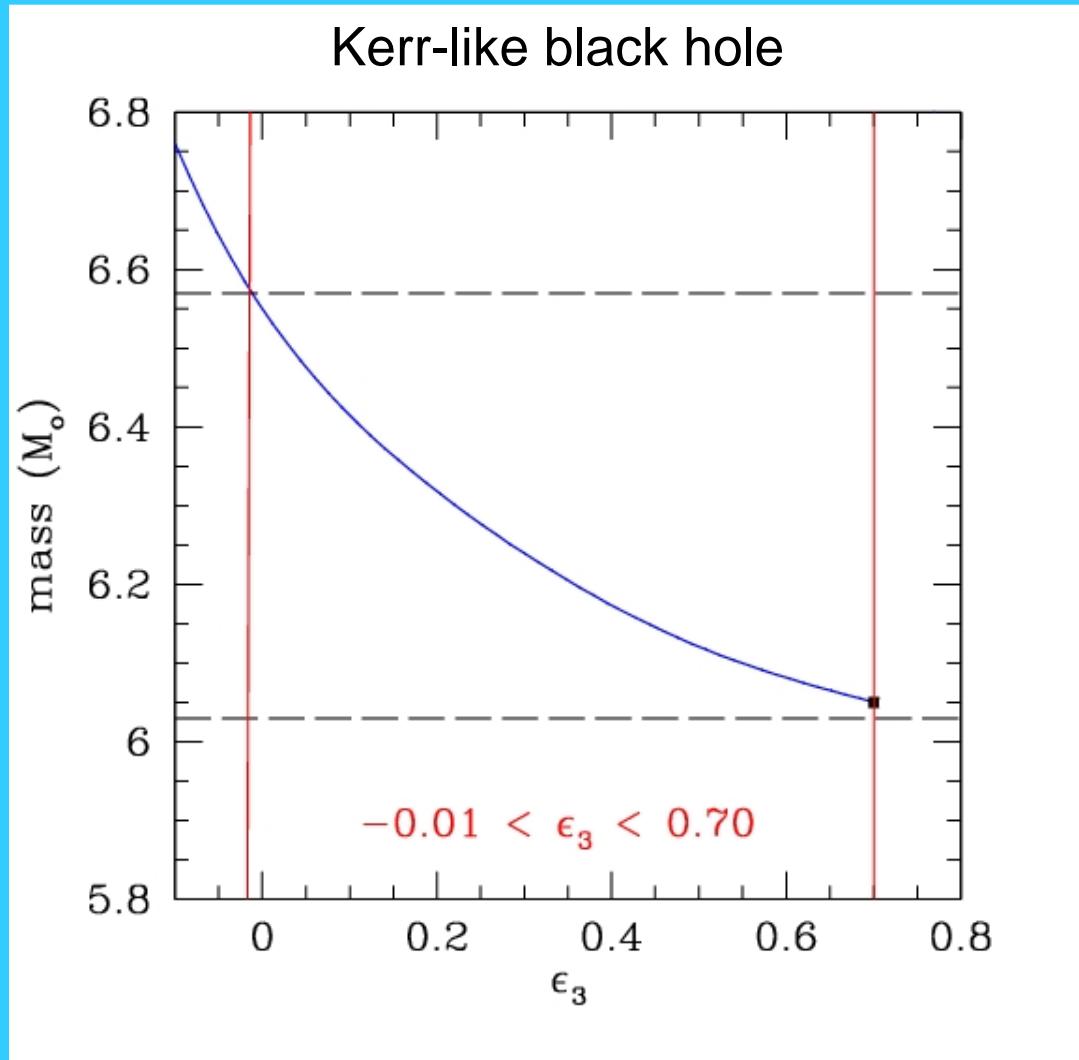


After Wagoner et al. 2001, ApJ, 559, L25
also Psaltis in X-Ray Timing 2003: Rossi and Beyond



After Wagoner et al. 2001, ApJ, 559, L25
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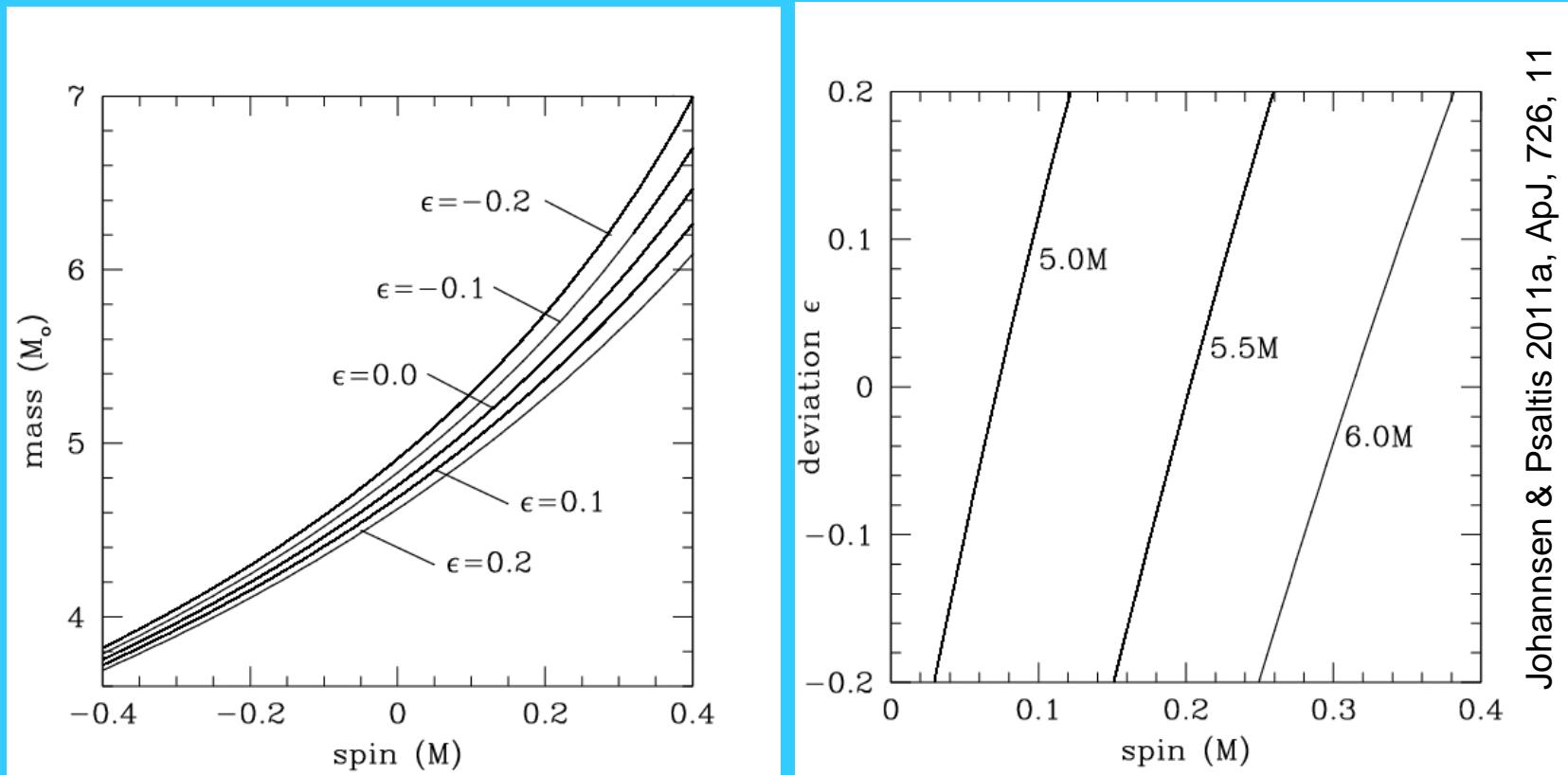




Test no-hair theorem with 2 QPOs + mass measurement !!

2. Resonance Model

Abramowicz & Kluzniak 2001, A&A, L19



1:2 resonance between Keplerian and radial epicyclic frequencies

Test no-hair theorem with 2 QPOs + mass + spin measurement

Summary and Outlook

- ✓ No-hair theorem characterizes nature of black holes
- ✓ Can be tested by measuring 3 multipole moments
- ✓ Observational tests: Imaging, Iron Lines, QPOs
- ✓ Diskoseismology model: 2 QPOs + mass measurement
- ✓ Resonance model: 2 QPOs + mass + spin measurement

Need Timing Mission to detect higher order harmonics and resolve model degeneracy

e.g., LOFT, arXiv:1107.0436